

A NEW PARAMETERS AND DIFFERENTIAL EQUATIONS OF ROTATION IN ORIENTATION OF A RIGID BODY PROBLEM

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Kinematics and dynamic differential equations of rotation of a rigid body for the «non-classical» conjugate three-dimensional new vectors of rotation are considered in this article. The vectors modules contain tangent and cotangent of one-fourth of the rotation angle. This paper also delivers applications of the conjugate equations in the dynamic and of orientation problems of a rigid body. Different new (polar) dynamical differential equations of RB rotation come out on the basis of kinematics equations.

In classical *Euler* occasion, for example, problem of decision of the system of six dynamical differential equations of *Euler-Poisson* reduces to integration of the system of just three dynamical equations with two independent classical first integrals of energy and square. Three new "tangent" vector coordinates-parameters definitely determine RB orientation in contrast to three direction cosines determined orientation of just one unit vector of supporting basis I in *Poisson* equations.

New vectors and conjugated polar equations define a (essentially) new line of fundamental research in solid mechanics.

Keywords: rotation vector *Rodrigues*, equations *Euler-Poisson*, orientation, rigid body, group rotation.

1. Three-dimensional vectors of rotation with minimal possible quantity of generalized coordinate lining with three freedoms of rigid body (RB)[1-4] with fixed point are of the most interest in the problem of dimensional orientation such as a task of orientation order parameter identification and directing of rotation (spherical) motion of RB.

Classical *Rodrigues* vector [1,4] has a module with tangent ($tg(\varphi/2)$) of half-angle φ of finite rotation of RB. This vector can not be used in the condition of $\varphi=\pi(180^\circ)$ in the problems of RB dimensional orientation. Considered in the works [2-4] tangent, cotangent *conjugate rotation vectors* $\tau=\tau k$, $\rho=\rho k$, where $\tau=k_\tau tg(\varphi/4)$, $\rho=k_\rho ctg(\varphi/4)$ (k_τ , k_ρ — arbitrary constant coefficients), k — unit vector of *Euler* axis of RB finite rotation are of the most interest (by $0<\varphi<2\pi$) because of its conjugate features [2,4,5].

2. The article deals with *conjugative polar* non-linear vector kinematics differential equations of RB rotation in the transformation of lineal vector $\omega(t)$ of RB angular velocity [3,4]

$$\boldsymbol{\tau}^* = \boldsymbol{\tau}' \boldsymbol{\Theta}_\tau \boldsymbol{\omega}, \quad \boldsymbol{\rho}^* = \boldsymbol{\rho}' \boldsymbol{\Theta}_\rho^T \boldsymbol{\omega}, \quad (1)$$

where $\boldsymbol{\tau}^* = d\boldsymbol{\tau}/dt$, $\boldsymbol{\rho}^* = d\boldsymbol{\rho}/dt$ - local derivative of vectors in the time t (derivative towards some connected with RB coordinate base); $\boldsymbol{\tau}' = \partial\boldsymbol{\tau}/\partial\varphi$, $\boldsymbol{\rho}' = \partial\boldsymbol{\rho}/\partial\varphi$ - particular derived modules $\boldsymbol{\tau}$, $\boldsymbol{\rho}$ of vectors of angle φ , defined as function

$$\boldsymbol{\tau}' = (k_\tau + \tau^2/k_\tau)/4, \quad \boldsymbol{\rho}' = -(k_\rho + \rho^2/k_\rho)/4; \quad \boldsymbol{\Theta}_\tau, \quad \boldsymbol{\Theta}_\rho^T$$

- orthogonal operators *semi-rotation* of RB (angle $\varphi/2$) [3]:

$$\begin{aligned} \boldsymbol{\Theta}_\tau &= \mathbf{E} + 2(k_\tau \mathbf{T} + \mathbf{T}^2)/(k_\tau^2 + \tau^2), \\ \boldsymbol{\Theta}_\rho^T &= \mathbf{E} + 2(-k_\rho \mathbf{R} + \mathbf{R}^2)/(k_\rho^2 + \rho^2), \end{aligned}$$

«T» - transposing.

Therein \mathbf{T} , \mathbf{R} - skew-symmetric operators of vector multiplication [2, 4], satisfies the identities $\mathbf{T}\boldsymbol{\tau} = \boldsymbol{\tau} \times \boldsymbol{\tau} = \mathbf{0}$, $\mathbf{R}\boldsymbol{\rho} = \boldsymbol{\rho} \times \boldsymbol{\rho} = \mathbf{0}$ - vector products, $\mathbf{0}$ - zero vector, \mathbf{E} - unit operator. Operators $\boldsymbol{\Theta}_\tau$, $\boldsymbol{\Theta}_\rho$ satisfy the identity $(\boldsymbol{\Theta}_\tau \boldsymbol{\Theta}_\rho)^2 = \mathbf{E}$, determined conjugacy of vectors $\boldsymbol{\tau}$, $\boldsymbol{\rho}$ and equations (1) [4] as features of “duality” and isomorphic correspondence.

3. Equations (1) have the first («trigonometric») integral in the form of scalar product $(\boldsymbol{\tau} \cdot \boldsymbol{\rho}) = \tau \rho = k_\tau k_\rho = C$ (arbitrary constant) and admit simple and graphic kinematics interpretation - polar precession-nutation model [4] with arbitrary vector $\boldsymbol{\omega}_\tau$ and for any RB. In the first equation vector $\boldsymbol{\omega}$ is transformed (at every moment of the traveling time of RB) by operator $\boldsymbol{\Theta}_\tau$ into vector $\boldsymbol{\omega}_\tau$ as a result of precession $\boldsymbol{\omega}_\tau$ on the angle $\psi_\tau = \varphi/2$ (rotation vector $\boldsymbol{\omega}$ with module $\omega(t)$ about the surface of some circular cone of «velocity» precession cone angle $2\nu_\omega$) around *Euler* axis with unit vector \mathbf{k} . Angle ν_ω - nutation angle (vector

deviation $\boldsymbol{\omega}$ from the unit vector \mathbf{k}). And then a vector $\boldsymbol{\omega}_\tau$ is multiplied by scalar operator $\boldsymbol{\tau}'\mathbf{E}$. At the second equation vector $\boldsymbol{\omega}$ precesses (rotates in the surface of that) at the angle $\psi_{\Theta\rho} = (\pi + \varphi/2)$ and is multiplied by scalar operator $\boldsymbol{\rho}'\mathbf{E}$. Angle ν_ω of nutation is determined from scalar product $(\boldsymbol{\omega} \cdot \boldsymbol{\tau})$, $(\boldsymbol{\tau}^* \cdot \boldsymbol{\tau})$ or $(\boldsymbol{\omega} \cdot \boldsymbol{\rho})$, $(\boldsymbol{\rho}^* \cdot \boldsymbol{\rho})$.

On the basis of such interpretation equation - models (1) come out, for example, kinematics anholonomic equality:

$$(\mathbf{a} \cdot \mathbf{b})^2 + ((\mathbf{a} \times \mathbf{b}) \cdot (\mathbf{a} \times \mathbf{b})) = (\mathbf{a} \cdot \mathbf{a})(\mathbf{b} \cdot \mathbf{b}),$$

$$\text{where } \mathbf{a} = \boldsymbol{\omega} \times \boldsymbol{\tau}, \quad \mathbf{b} = (\boldsymbol{\Theta}_\tau \boldsymbol{\omega}) \times \boldsymbol{\tau}.$$

Equations (1) have general decisions in the form of *Cauchy* [4, 5]. General decision of the first equation in (1) determines the form shape of the new vector representations of three-dimensional rotation (a *Lie* group) [6]. In this group inverse element $\boldsymbol{\tau}^{-1}$ is equal to adverse vector, i. e. $\boldsymbol{\tau}^{-1} = -\boldsymbol{\tau}$, but unit element is equal to zero vector $\mathbf{0}$.

4. Different new (polar) dynamical differential equations of RB rotation come out on the basis of kinematics equations (1).

In classical *Euler* occasion, for example, problem of decision of the system of six dynamical differential equations of *Euler-Poisson* [7] reduces to integration of the system of just three dynamical equations with two independent classical first integrals (of energy and square [7]). These two integrals are enough to consider the system of three equations to be integrated [7].

At this case change is introduced in the equations (1) $\boldsymbol{\omega} = \mathbf{S}^{-1} \mathbf{g}$, where \mathbf{g} - constant vector of kinetic moment (constant module and in the line of supporting basis I [7]); \mathbf{S}^{-1} - inverse

operator S of RB inertia [3,4,7] (constant in connected with RB basis J).

Contained of (1), for example, polar matrix differential equation with vector coordinates-parameters τ has the form (see also [3]):

$$\tau^*_{\tau J} = \tau' (\Theta_{\tau J} S_J^{-1} \Theta_{\tau J}^T) \Theta_{\tau J}^T g_I, \quad (2)$$

where $\tau^*_{\tau J} = [\dot{\tau}_{j1} \dot{\tau}_{j2} \dot{\tau}_{j3}]^T$ - column matrix with derivative coordinates τ^* in basis J ; $\Theta_{\tau J}$ - matrix (3x3) of operator Θ_{τ} in basis J ; S_J - diagonal matrix (3x3) with three constant the main moments of RB inertia; S_J^{-1} - inverse matrix, $g_I = [g_{i1} g_{i2} g_{i3}]^T$ - constant column matrix with vector coordinates g in basis I .

Three coordinates of the new vector-parameter τ is uniquely determined RB orientation (as opposed to the three direction cosines of only one unit ortho-vertical supporting basis I in the *Poisson* equations [7]). In the problems of the dynamics of RB and synthesis problems of attitude control laws (definition of control moments) with the use of *Lyapunov* functions (see, eg, [8]), the scalar equations (1), (2) are used in conjunction with the classical dynamic *Euler* equations [1, 7, 8].

5. Conjugate vectors τ, ρ are considering as parts of vectors with non-traditional non-normalized (non-*Hamiltonian*) rotation new quaternions [8]. Such non-normalized quaternions can be effectively used instead of classical normalized quaternions (with parameters of *Rodrigues-Hamilton, Euler-Rodrigues* [4]) in different «particular» tasks of

inertial orientation [1, 4, 9, 10], including RB orientation control problem [8].

Sets of coordinates-parameters of conjugate vectors τ, ρ , corresponded to them quaternions sufficiently broaden the circle of new RB orientation parameters. Such parameters can be of interest in the problem of global parameterization [9] of three-dimensional rotation group and new representation *Lorenz* group in quantum mechanics [11].

New vectors τ, ρ and conjugated polar equations (1), (2) define the (essentially) the new direction of fundamental research in solid mechanics.

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